

LETTER

## Ion microprobe analysis of $(^{231}\text{Pa})/(^{235}\text{U})$ and an appraisal of protactinium partitioning in igneous zircon

AXEL K. SCHMITT\*

Department of Earth and Space Sciences, University of California, Los Angeles, California 90095, U.S.A.

### ABSTRACT

Ion microprobe U-Pa measurements in zircon crystals from a Holocene rhyolite (Salton Buttes, California) are utilized to assess the partitioning of Pa and U between zircon and melt. The relative sensitivity factor (RSF) for ion microprobe U-Pa zircon analysis was calibrated on a natural high-U secular equilibrium zircon from Buff Peak (Nevada). Zircon crystals from Salton Buttes rhyolite show evidence for excess  $^{231}\text{Pa}$  with a weighted average  $(^{231}\text{Pa})/(^{235}\text{U})$  activity ratio of 1.9. From this ratio, model zircon-melt partition coefficient ratios  $D_{\text{Pa}}/D_{\text{U}} = 0.9\text{--}2.2$  are obtained, assuming plausible limits for  $(^{231}\text{Pa})/(^{235}\text{U})_{\text{melt}}$  and correcting for the age of zircon crystallization ( $\sim 15$  ka). These values roughly fit lattice strain models for tetravalent cations, but are more than one order of magnitude lower than predictions for pentavalent Pa. Based on this appraisal of Pa partitioning, initial  $^{231}\text{Pa}$  disequilibrium caused by zircon-melt fractionation alone can produce minor discordance or excess  $^{207}\text{Pb}$  in Late Cenozoic zircon, but is insignificant for older U-Pb zircon ages.

**Keywords:** Pa-231, zircon, partition coefficients, uranium disequilibrium, rhyolites

### INTRODUCTION

Zircon is of paramount importance for geochronology based on its preference for actinide elements while excluding initial Pb, its chemical and physical stability, and slow diffusion of radioactive parents and daughters. It is amenable to multiple dating methods, including U-Pb, Th-Pb, U-Th-He, radiation damage, and U-series disequilibrium. Disequilibrium between long-lived intermediate daughters  $^{230}\text{Th}$  (half-life  $t_{1/2} = 75\,690$  a) or  $^{231}\text{Pa}$  ( $t_{1/2} = 32\,760$  a) and their respective parents  $^{238}\text{U}$  (magmatic fractionation between  $^{234}\text{U}$  and  $^{238}\text{U}$  assumed negligible) and  $^{235}\text{U}$  ultimately impacts the abundance of radiogenic  $^{206}\text{Pb}$ ,  $^{207}\text{Pb}$ , and He, and requires corrections for accurate age determination (e.g., Mattinson 1973). Zircon favors U over Th with a zircon-melt partitioning coefficient ratio  $D_{\text{U}}/D_{\text{Th}} \sim 6$  (e.g., Blundy and Wood 2003) so that initial  $(^{230}\text{Th})/(^{238}\text{U})_{\text{zircon}}$  (activities denoted by parentheses) is expected to be  $\sim 0.17$ . The resulting  $^{206}\text{Pb}$  deficit is practically negligible within analytical precision for zircon crystals older than  $\sim 100$  Ma and even for complete exclusion of initial  $^{230}\text{Th}$  can only amount to a  $\sim 108$  ka age difference (Schärer 1984). By contrast, excess  $^{231}\text{Pa}$  has been frequently invoked as a source for strong  $^{207}\text{Pb}$  excesses in Pre-Quaternary or even Archean zircon crystals, including those of the oldest known terrestrial zircon (e.g., Anczkiewicz et al. 2001; Parrish and Noble 2003; cf. Wilde et al. 2001). In theory, the incorporation of strong  $^{231}\text{Pa}$  excesses in zircon could significantly affect even old  $^{207}\text{Pb}/^{235}\text{U}$  and  $^{207}\text{Pb}/^{206}\text{Pb}$  ages.

Lattice strain models for zircon-melt partitioning of  $\text{Pa}^{5+}$  (Fig. 1; Blundy and Wood 2003) are based on estimates for  $\text{Nb}^{5+}$  as a proxy for  $\text{Pa}^{5+}$ . Published  $D_{\text{Nb}} = 0.34$  and lattice strain parameters (apparent Young's modulus  $E^{5+} > E^{4+} = 750$  GPa and site

radius  $r_0^{4+} = r_0^{5+} = 0.91$  Å; Blundy and Wood 2003) imply that the parabola for  $5+$  cations should be similar to the  $4+$  curve, and consequently  $\text{Pa}^{5+}$  is expected to be highly compatible in zircon ( $D_{\text{Pa}}/D_{\text{U}} > 30$ ; Fig. 1). Direct experimental constraints on zircon-melt partitioning of Pa, however, are lacking. The low abundance of  $^{231}\text{Pa}$ , high radioactivity and toxicity, and the absence of a stable Pa isotope impose such analytical and experimental challenges that Parrish and Noble (2003) commented that the  $^{231}\text{Pa}$  isotope effects in accessory minerals “can be inferred only by the measurements of the decay products of  $^{231}\text{Pa}$ , not to the measurement of  $^{231}\text{Pa}$  directly” (p. 205). Here, I introduce a high-sensitivity and high-spatial resolution secondary ionization mass spectrometry (SIMS, ion microprobe) method for the direct determination of  $^{231}\text{Pa}$  in zircon. By applying this technique to young zircon crystals from a rhyolite with a Holocene eruption age, the first observational constraints on zircon-melt partitioning of Pa relative to U are obtained.

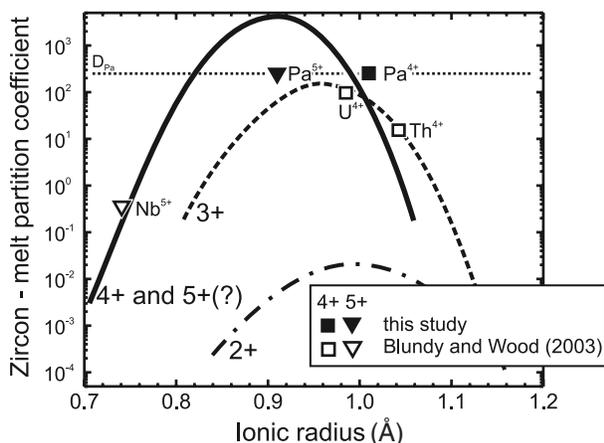
### METHODS

Hand-picked zircon grains (typically 100–150  $\mu\text{m}$  long) were epoxy-mounted, ground, and polished with SiC paper and 1  $\mu\text{m}$  diamond paste, and Au-coated. Ion microprobe analysis on the UCLA Cameca ims 1270 used a mass-filtered  $\text{O}^+$  ( $\text{O}$  refers to  $^{16}\text{O}$  unless specifically noted) primary ion beam at  $\sim 60$  nA and 22.5 keV impact energy, focused to a  $\sim 30 \times 40$   $\mu\text{m}$  spot. Secondary ions were extracted at 10 kV with an energy bandpass of 50 eV and detected with a high-resistivity ( $10^{11}$   $\Omega$ ) Faraday cup (FC) counting system for intensities  $> 10^6$  counts per second (cps) (typically for  $^{232}\text{ThO}^+$  and  $^{238}\text{UO}^+$ ) and with an ETP electron multiplier (EM) operated at  $\sim 2$  kV for ion intensities  $< 10^6$  cps for all other masses. Count rates were corrected for EM dead time (25 ns) and FC baseline recorded on mass 247.5 during each measurement ( $\sim 3 \times 10^5$  cps). FC-EM relative yields between 1.01 and 1.02 were determined from the ratio of measured  $^{238}\text{UO}^+/^{235}\text{UO}^+$  to the  $^{238}\text{U}/^{235}\text{U}$  atomic abundance ratio of 137.88 and used for correction of  $^{230}\text{ThO}^+/^{232}\text{ThO}^+$  ratios.

$^{231}\text{PaO}^+$  molecular ions are  $\sim 10$  times more abundant than  $^{231}\text{Pa}^+$ , similar to the behavior of U and Th atomic and molecular species in SIMS positive ion analysis of zircon. High-resolution mass spectra around mass 246 and 247 on natural zircons, synthetic cubic zirconia, hafnia, and the epoxy mounting medium indicate that a

\* E-mail: axel@argon.ess.ucla.edu

mass resolving power  $m/\Delta m = \sim 4000$  is sufficient to resolve the  $^{197}\text{Au}^{18}\text{OO}_2^+$  and  $^{197}\text{AuO}_3\text{H}_2^+$  interferences from  $^{231}\text{PaO}^+$  (Fig. 2). Unresolved isobaric interferences such as  $^{230}\text{Th}^{17}\text{O}^+$  and  $^{230}\text{ThOH}^+$  are insignificant due to low abundances of the calculated peaks based upon the 246 amu mass region. Zirconium oxide molecular ion peaks (Fig. 2) were used for calibrating the magnetic field settings for low abundance isotopes and background mass stations at the beginning of each analysis



**FIGURE 1.** Lattice strain model for zircon-melt partitioning with curves for 4+ (solid), 3+ (dash), and 2+ (dash-dot) cations (after Blundy and Wood 2003). Model  $D_{\text{Pa}}$  (this study; uncertainty equivalent to size of symbols—see text) is plotted against ionic radius for  $\text{Pa}^{4+}$  and  $\text{Pa}^{5+}$  (VIII-coordination) using  $D_{\text{U}} = 100$ . Tetravalent  $D_{\text{Th}}$  and  $D_{\text{Pa}}$  both slightly deviate from the 4+ curve, whereas  $D_{\text{Pa}}$  for  $\text{Pa}^{5+}$  is more than 10× lower than predicted from using  $D_{\text{Nb}}$  as a proxy.

and after every 15 cycles; on more intense peaks the magnetic field settings were calibrated directly in the same intervals. The abundance sensitivity at mass 247 due to the  $^{232}\text{ThO}^+$  peak ( $\sim 250$  ppb measured on synthetic thorite) contributes to the background at mass 247.03 and intensities recorded at 247.1 amu were subtracted from the  $^{231}\text{PaO}^+$  peak (Table 1). At a sputter rate of  $\sim 0.05 \mu\text{m}^3/\text{s/nA O}^-$ , approximately 30 ng of zircon or 0.3 fg of Pa are consumed during a typical 30 min analysis. This corresponds to a useful yield (atoms detected over atoms sputtered) for  $^{231}\text{PaO}^+$  in zircon of  $\sim 0.92\%$ , similar to the useful yield determined for  $^{230}\text{ThO}^+$  ( $\sim 1.04\%$ ).

## RESULTS

### Relative sensitivity calibration on Buff Peak zircon (BP-1)

Relative sensitivity for Th/U species in SIMS is commonly calibrated by analyzing radiogenic  $^{208}\text{Pb}/^{206}\text{Pb}$  on a concordant zircon with known age (Reid et al. 1997). Because a separate Pa-Pb decay system is non-existent and Pa-zircon standards are unavailable, the instrumental mass fractionation was determined from measuring the ratio of  $^{231}\text{PaO}^+ \times \lambda_{231}$  over  $^{235}\text{UO}^+ \times \lambda_{235}$  on a secular equilibrium reference zircon with  $(^{231}\text{Pa})/(^{235}\text{U}) = 1$  ( $\lambda$  refers to decay constant). Reference zircons were extracted from the 14.57 Ma ( $^{40}\text{Ar}/^{39}\text{Ar}$  sanidine single-crystal age) Buff Peak tuff (Nevada), a rhyolite with  $\sim 20$  ppm U, which is among the highest bulk rock U abundances for volcanic rocks in the Western U.S.A. (Castor and Henry 2000). Buff Peak zircons (BP-1) are homogenous in back-scatter electron imaging, except for occasionally intergrown thorite domains, which were avoided. Their high U abundance ( $\sim 1\text{--}4$  wt%  $\text{UO}_2$ ) and comparatively high U/Th ( $\sim 2.7\text{--}4.8$ ) are favorable for  $^{231}\text{Pa}$  analysis (Table 1).

**TABLE 1.** Ion microprobe U-series results for secular equilibrium reference zircons (BP-1, AS-3) and Salton Buttes rhyolite zircons

Zircon	grain*	spot	$(^{230}\text{Th})/(^{238}\text{U})$	$(^{238}\text{U})/(^{232}\text{Th})$	$(^{230}\text{Th})/(^{232}\text{Th})$	$(^{231}\text{Pa})/(^{235}\text{U})$	246.038/246.3	247.031/247.1	247.031 (cps)	$\text{UO}_2$ (wt%)
secular equilibrium (wt%)										
BP-1†	m1g2	1	$1.012 \pm 0.012$	$8.233 \pm 0.082$	$8.334 \pm 0.051$	—	617	—	—	$3.08 \pm$
BP-1	m1g2	2	$1.019 \pm 0.014$	$8.470 \pm 0.086$	$8.631 \pm 0.086$	$1.047 \pm 0.074$	591	4.60	5.21	$3.23 \pm$
BP-1	m1g3	1	$0.985 \pm 0.015$	$11.13 \pm 0.12$	$10.97 \pm 0.11$	$1.042 \pm 0.067$	1112	5.36	5.36	$1.77 \pm$
BP-1	m1g4	1	$0.985 \pm 0.017$	$14.66 \pm 0.15$	$14.43 \pm 0.20$	$0.916 \pm 0.075$	859	5.41	2.41	$0.87 \pm$
BP-1	m1g5	1	$0.974 \pm 0.018$	$11.72 \pm 0.13$	$11.41 \pm 0.16$	$1.012 \pm 0.068$	805	5.18	4.89	$1.64 \pm$
BP-1	m1g6	1	$0.993 \pm 0.013$	$12.60 \pm 0.13$	$12.51 \pm 0.10$	$1.098 \pm 0.070$	1316	5.48	5.22	$1.49 \pm$
BP-1	m1g7	1	$0.976 \pm 0.022$	$10.09 \pm 0.18$	$9.850 \pm 0.133$	$0.894 \pm 0.062$	1566	3.19	6.16	$1.79 \pm$
BP-1	m2g1	1	$0.985 \pm 0.011$	$9.389 \pm 0.076$	$9.248 \pm 0.065$	$0.972 \pm 0.065$	1039	3.58	7.09	$3.07$
BP-1	m2g2	1	$0.976 \pm 0.011$	$12.10 \pm 0.10$	$11.81 \pm 0.09$	$1.029 \pm 0.065$	818	3.96	6.32	$3.96$
BP-1	m2g3	1	$0.978 \pm 0.010$	$10.23 \pm 0.09$	$10.00 \pm 0.06$	$0.905 \pm 0.062$	485	3.77	7.48	$3.35$
BP-1	m2g4	1	$1.017 \pm 0.010$	$9.142 \pm 0.071$	$9.298 \pm 0.06$	$1.110 \pm 0.068$	593	4.59	8.34	$2.99$
										U (ppm)
AS-3	m1g2	1	$0.964 \pm 0.017$	$4.532 \pm 0.045$	$4.369 \pm 0.064$	$1.113 \pm 0.166$	124	2.02	0.36	490
Salton Buttes										
SB0402§	m1g3	1	$0.313 \pm 0.024$	$4.783 \pm 0.077$	$1.497 \pm 0.112$	$3.281 \pm 0.672$	36.8	2.51	0.40	1150
SB0402	m1g1	1	$0.237 \pm 0.022$	$5.833 \pm 0.065$	$1.385 \pm 0.130$	$2.008 \pm 0.535$	14.8	1.98	0.29	240
SB0402	m1g4	1	$0.269 \pm 0.015$	$5.483 \pm 0.113$	$1.475 \pm 0.079$	$2.011 \pm 0.375$	32.6	4.56	0.45	570
SB0401	m2g1	1	$0.272 \pm 0.012$	$4.865 \pm 0.048$	$1.324 \pm 0.057$	$2.180 \pm 0.337$	272	4.14	0.85	1650
SB0401	m2g3	1	$0.320 \pm 0.017$	$5.011 \pm 0.051$	$1.605 \pm 0.083$	$2.344 \pm 0.489$	22.2	3.79	0.33	490
SB0402	m2g1	1	$0.262 \pm 0.018$	$5.412 \pm 0.055$	$1.420 \pm 0.096$	$1.802 \pm 0.465$	31.2	2.40	0.34	590
SB0402	m2g2	1	$0.213 \pm 0.025$	$7.226 \pm 0.099$	$1.538 \pm 0.182$	$3.776 \pm 1.139$	15.5	6.09	0.30	340
SB0402	m2g4	1	$0.258 \pm 0.021$	$6.241 \pm 0.050$	$1.610 \pm 0.133$	$2.562 \pm 0.878$	28.5	3.21	0.29	430
SB0401	m2g4	1	$0.273 \pm 0.014$	$5.156 \pm 0.093$	$1.407 \pm 0.067$	$1.787 \pm 0.706$	37.2	4.00	0.32	580
SB0401	m2g2	1	$0.395 \pm 0.010$	$2.834 \pm 0.032$	$1.120 \pm 0.025$	$1.536 \pm 0.186$	77.1	2.45	1.66	2940

Note: Activities in parentheses; decay constants used:  $\lambda_{235} = 9.8485 \times 10^{-10} \text{ a}^{-1}$ ;  $\lambda_{231} = 2.1158 \times 10^{-5} \text{ a}^{-1}$ ;  $\lambda_{238} = 1.55125 \times 10^{-10} \text{ a}^{-1}$ ;  $\lambda_{230} = 9.1580 \times 10^{-6} \text{ a}^{-1}$ ;  $\lambda_{232} = 4.9475 \times 10^{-11} \text{ a}^{-1}$  (compiled in Bourdon et al. 2003). All errors  $1\sigma$  (including measurement and RSF uncertainties) 246.038: peak  $^{230}\text{ThO}^+$ ; 246.3: background  $^{230}\text{ThO}^+$ ; 247.031: peak  $^{231}\text{PaO}^+$ ; 247.1 background  $^{231}\text{PaO}^+$ .

\* m1 analyzed August 22, 2006, (RSF Th/U = 0.908±0.009; RSF Pa/U = 0.782±0.020; RSF U/Zr = 0.02); m2 analyzed August 30, 2006, (RSF Th/U = 0.908±0.007; RSF Pa/U = 0.783±0.023; RSF U/Zr = 0.014); relative sensitivity factor (RSF) Th/U from  $^{208}\text{Pb}/^{206}\text{Pb}$  analysis of 91500 zircon; RSF Pa/U from BP-1;  $(^{231}\text{Pa})/(^{235}\text{U})$  of BP-1 only shown as indication of reproducibility U/Zr RSF from  $^{238}\text{UO}^+ / ^{90}\text{ZrO}_2$  on 91500 (81.2 ppm U).

† Sample location: N 41°21'37.1"; W 118°19'21.2".

‡ Electron microprobe analysis (in wt%):  $\text{SiO}_2 = 30.4\text{--}31.1$ ;  $\text{P}_2\text{O}_5 = 0.23\text{--}0.62$ ;  $\text{Y}_2\text{O}_3 = 1.45\text{--}2.65$ ;  $\text{ZrO}_2 = 57.4\text{--}61.3$ ;  $\text{HfO}_2 = 1.59\text{--}1.90$ ;  $\text{ThO}_2 = 0.29\text{--}1.04$ ;  $\text{UO}_2 = 1.36\text{--}3.64$ ; totals = 97.2–99.4.

§ SB0402 sampled at Obsidian Butte: N 33°10'14.6"; W 115°38'03.1".

|| SB0401 sampled at Red Island: N 33°11'59.1"; W 115°36'44.3".



obtained for Salton Buttes zircon (Figs. 1 and 4).

The resulting model zircon-melt  $D_{Pa}$  for  $Pa^{5+}$  is much lower than predicted from its proxy  $D_{Nb}$  (Fig. 1; Blundy and Wood 2003). It is conceivable that this discrepancy is caused by an over-determination of Nb in zircon due to zirconium hydride interference (Blundy and Wood 2003). Interestingly, model  $D_{Pa}$  for  $Pa^{4+}$  plots closely to but not perfectly on the lattice strain curve for  $4+$  cations, leaving the possibility that  $Pa^{5+}$  species were insignificant in Salton Buttes rhyolite. Anomalously reduced magmatic conditions, however, can be ruled out based on characteristic positive Ce anomalies of Salton Buttes zircon ( $Ce/Ce^* = 26-60$ ; Schmitt and Vazquez 2006).

Regardless of Pa speciation, the detection of  $(^{231}Pa)/(^{235}U)$  disequilibrium in young zircon implies that Pa-U dating has the potential to establish relative crystallization sequences for late Pleistocene-Holocene (<150 ka) zircon if initial  $(^{231}Pa)/(^{235}U)$  was uniform for a suite of crystals in a rock, or yield absolute ages if the initial  $(^{231}Pa)/(^{235}U)$  can be independently constrained. The conclusion for U-Pb geochronology is that differences between concordant  $^{207}Pb$  and excess  $^{207}Pb$  ( $\Delta^{207}Pb$ ; Schärer 1984) caused by zircon-melt fractionation with  $D_{Pa}/D_U \sim 0.9-2.2$  are vanishingly small for zircon older than a few ~tens Ma. Previous studies invoked much higher  $\Delta^{207}Pb$  resulting from  $^{231}Pa$  disequilibrium in zircon (e.g., ~100% for ca. 30 Ma zircons and ~5% for Archean zircons; Anczkiewicz et al. 2001; Parrish and Noble 2003). This would require initial  $^{231}Pa$  excesses in the order of  $\sim 10^5-10^{70}\%$ , instead of the maximum of ~200 % constrained by model  $D_{Pa}/D_U$  (Fig. 1).  $^{207}Pb$  excesses in zircon—if valid—are therefore unlikely to result from zircon-melt fractionation alone and would require zircon crystallization in unrecognized reservoirs with extremely high  $(^{231}Pa)/(^{235}U)$ .

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